

## Measurement of Neutron Yield Produced by 135 MeV/nucleon $H_2$ and $^2D$ Beams Incident on a Thick Iron Target with the Activation Method

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### Abstract

The neutron production by 135 MeV/nucleon  $H_2$  and  $^2D$  beams incident on a thick iron target was measured with the activation method. Seven metals of C, Al, Fe, Co, Ni, In and Au were used as the detectors. Cross sections of neutron - induced reactions employed are assumed to be constant in the neutron energy region of more than several tens MeV for lack of data. The preliminary neutron yield are presented.

### I. INTRODUCTION

It is indispensable to have accurate data for neutron production reactions in designing shields

of accelerator, reactor, medical facilities and so on. Accumuration of those data, however, is extremely poor for intermediate energy accelerators. Actually, we were forced to take a convenient way in shielding calculation for the RIKEN Ring Cyclotron Facility.<sup>1)</sup> A series of measurements of neutron yields has been planed for reactions of various incident particles on a thick target of iron with activation detectors.<sup>2)</sup> Here, neutron spectra produced by 135 MeV/nucleon  $He_2$  and  $^2D$  incident on an iron target are given. Extensive analysis and experiments are in progress.

### II. EXPERIMENTAL

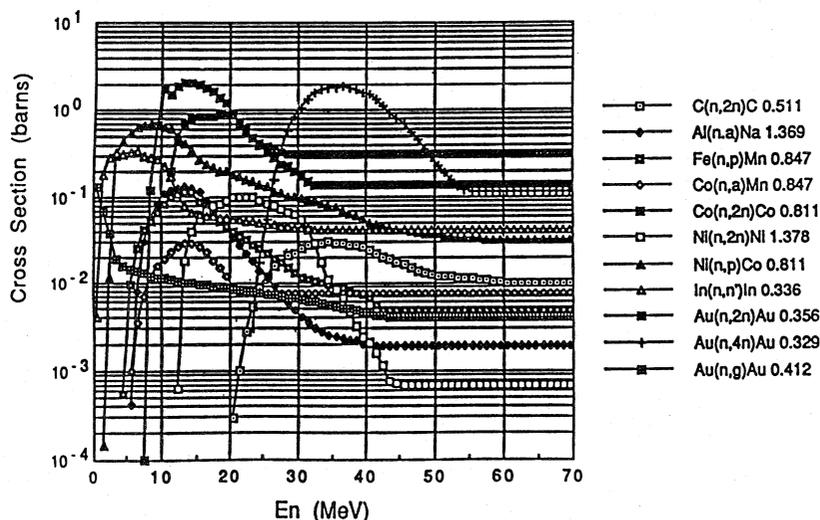


Fig.1 Cross section curves of employed reactions

The activation method has been used widely for the spectral measurement of neutrons produced by nuclear reactions. In many cases, however, the energies of incident particles were not so high, that is, less than 70 MeV, and proton beams only were used. A series of the experiments has made extension of kinds and energies of the incident particles. Having distinctive features, this method has two defects; (1) data of either activation cross sections or neutron fluxes are necessary because experiments give product of the two only, (2) some ambiguity in decision of one of them exists in unfolding process. We employ cross section data cited from Refs. 3~5. The used cross section curves are shown in Fig. 1. In the energy region of more than 70 MeV these cross sections are extrapolated to be constant.

Activation rates  $A_i$  [ $\text{cm}^2 \cdot \text{sr}^{-1} \cdot \text{particle}^{-1}$ ] is given by

$$A_i = \sum_j \sigma_i(E_j) \cdot \Phi(E_j)$$

$$= \frac{r^2 \cdot \lambda_i \cdot C_i \cdot \exp(-\lambda_i t_w)^{-1} \cdot \{1 - \exp(-\lambda_i t_c)\}^{-1}}{\epsilon_i \cdot \eta_i \cdot N_0 \cdot P_0 \cdot \{1 - \exp(-\lambda_i t_0)\}}$$

where

$\sigma_i$  = cross section of the  $i$ th isotope;

$E_j$  =  $j$ th energy segment;

$\Phi(E_j)$  = number of emitted neutrons per sr per incident particle;

$r$  = distance between the detector and the target Fe;

$\lambda_i$  = decay constant of the  $i$ th isotope;

$C_i$  = peak count rate;

$\epsilon_i$  = counting efficiency of the  $\gamma$  ray;

$\eta_i$  = branching ratio of the  $\gamma$  ray;

$N_0$  = number of nuclei in the detector;

$P_0$  = number of the incident particles per sec.;

$t_w$  = waiting time;

$t_c$  = counting time;

$t_0$  = irradiation time.

Table 1 Characteristics of activation detectors

Detector	Reaction	Isotopic abundance (%)	Half life	Decay const $\lambda$ ( $\text{sec}^{-1}$ )	$\gamma$ -ray energy (MeV)	Branching ratio
C	$^{12}\text{C}(n,2n)^{11}\text{C}$	98.89	20.38m	$5.669 \times 10^{-4}$	0.511	1.9952
Al	$^{27}\text{Al}(n,\alpha)^{24}\text{Na}$	100.0	15.02h	$1.282 \times 10^{-5}$	1.369	1.0000
Fe	$^{56}\text{Fe}(n,p)^{56}\text{Mn}$	91.68	2.579h	$7.467 \times 10^{-5}$	0.847	0.9887
Co	$^{59}\text{Co}(n,\alpha)^{56}\text{Mn}$	100.0	2.518h	$7.467 \times 10^{-5}$	0.847	0.9887
Co	$^{59}\text{Co}(n,2n)^{58}\text{Co}$	100.0	70.79d	$1.133 \times 10^{-7}$	0.811	0.9944
Ni	$^{58}\text{Ni}(n,2n)^{57}\text{Ni}$	68.27	35.94h	$5.357 \times 10^{-6}$	1.378	0.7768
Ni	$^{58}\text{Ni}(n,p)^{58}\text{Co}$	68.27	70.79d	$1.133 \times 10^{-7}$	0.811	0.9944
In	$^{115}\text{In}(n,n)^{115m}\text{In}$	95.67	4.486h	$4.292 \times 10^{-5}$	0.336	0.4670
Au	$^{197}\text{Au}(n,2n)^{196}\text{Au}$	100.0	6.183d	$1.298 \times 10^{-6}$	0.356	0.8770
Au	$^{197}\text{Au}(n,4n)^{194}\text{Au}$	100.0	3.955h	$4.868 \times 10^{-6}$	0.329	0.6130
Au	$^{197}\text{Au}(n,\gamma)^{198}\text{Au}$	100.0	2.696d	$2.976 \times 10^{-6}$	0.412	0.9550

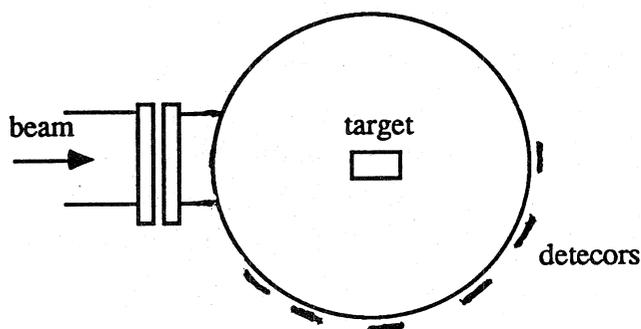


Fig. 2 experimental arrangement

Details of the activation detectors are given in Table 1. An irradiation chamber is shown schematically in Fig.2.

Unfolded spectra were obtained by using a small code based on the same algorithm as SAND-II.

### III RESULTS and DISCUSSIONS

Figure 3 shows angular dependence of activities of each detector by proton and deuteron beams incident on a thick iron target, respectively. Some rough features which will be seen from the figure are as follows: (1) Neutron yields are more in the case of deuteron than proton incidences. (2)

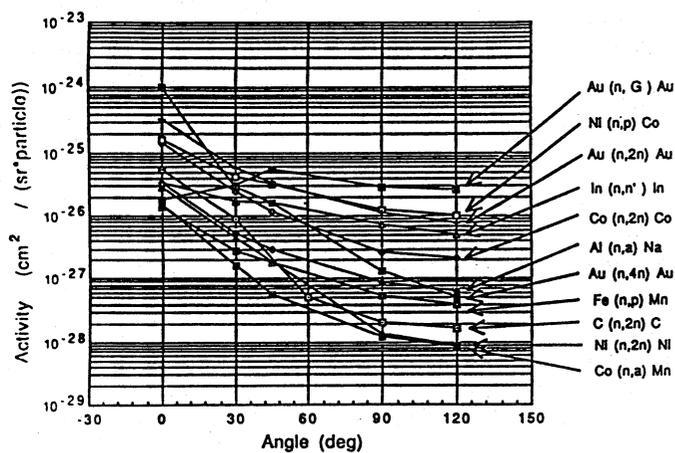
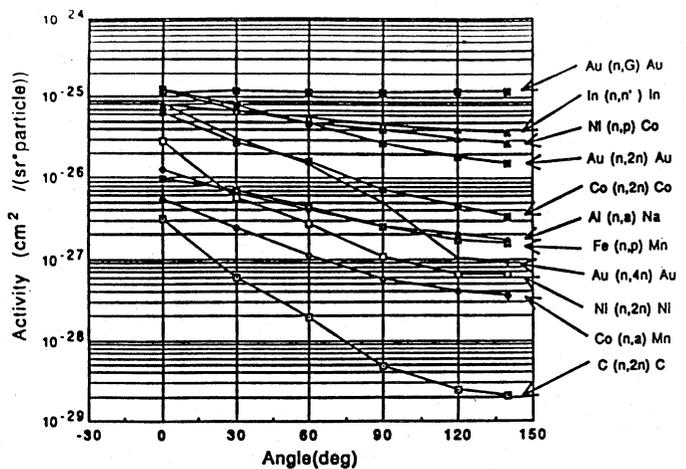


Fig.3 Angular dependence of activities for proton(upper) and deuteron(lower) incidences

Angular dependence of neutron yields is larger in reactions by deuterons than protons.

Preliminary neutron fluxes are shown in Fig.4. The activation method has essentially the defect that results are influenced by the first guess in unfolding process. A function  $\phi = \exp(-E_n/a)$  was used as first guesses, where  $a$ 's are taken between 5 and 20. Values and energy dependence of neutron fluxes were affected strongly by the values. It will be seen that low energy neutrons are rather more in proton reactions. Extensive analyses are in progress.

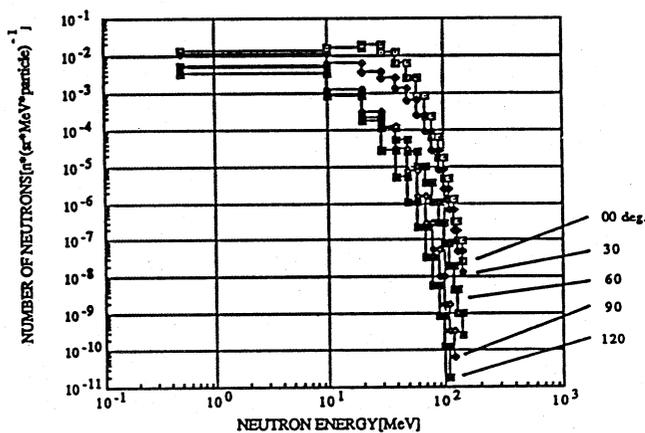
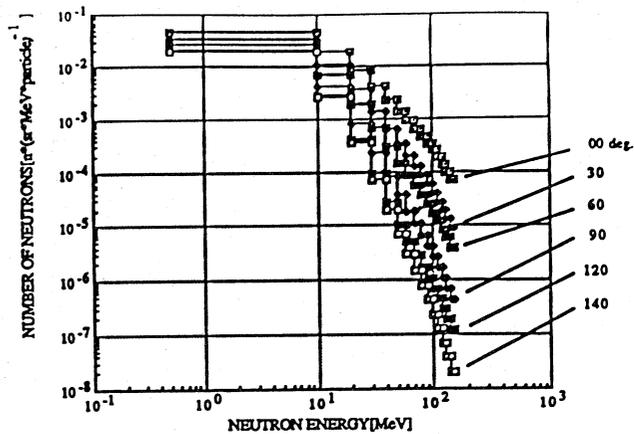


Fig.4 Preliminary neutron fluxes for proton (upper) and deuteron(lower) incidences

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